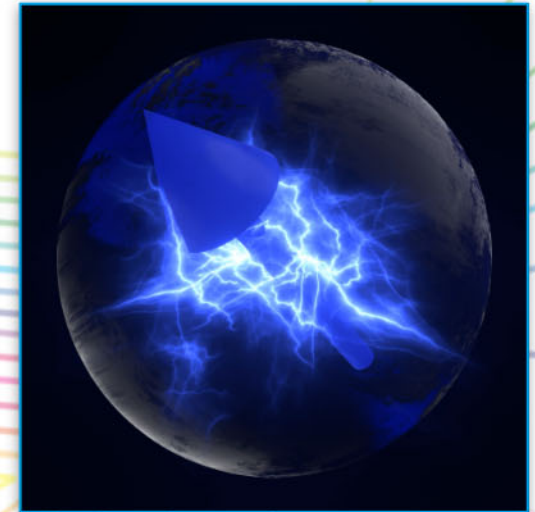


Spherical Neutron Polarimetry

A pioneering ILL technique and its unique contribution to multiferroics research

Navid Qureshi



Outline

Content of the presentation



Introduction to SNP



The **Cryogenic Polarisation Analysis Device**



Key contributions to magnetoelectric and multiferroic systems



Conclusions



Introduction to SNP

Motivation

THE PROBLEM

Unpolarised neutrons only give access to the modulus of the magnetic interaction vector \mathbf{M}_\perp via the observed intensity $I \sim \mathbf{M}_\perp \mathbf{M}_\perp^* = |\mathbf{M}_\perp|^2$.

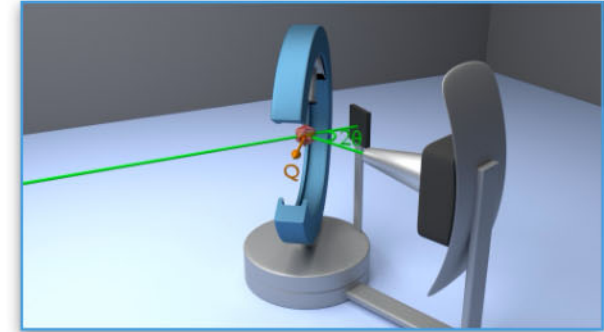
$$\mathbf{M}_\perp(\mathbf{Q}) = \hat{\mathbf{Q}} \times (\mathbf{M} \times \hat{\mathbf{Q}}) \quad \mathbf{M}(\mathbf{Q}) = \sum_j \mathbf{S}_j f_j(\mathbf{Q}) \exp(i\mathbf{Q}\mathbf{r}_j)$$

THE RAY OF HOPE

In 1963 Blume and Maleev independently publish their equations relating the final neutron polarisation to generalised cross sections.

THE DREAM

Be able to measure all the terms in these equations precisely.



Blume (1963) Phys. Rev. **130** 1670

Maleev et al. (1963) Sov. Phys. Solid State. **4** 2533

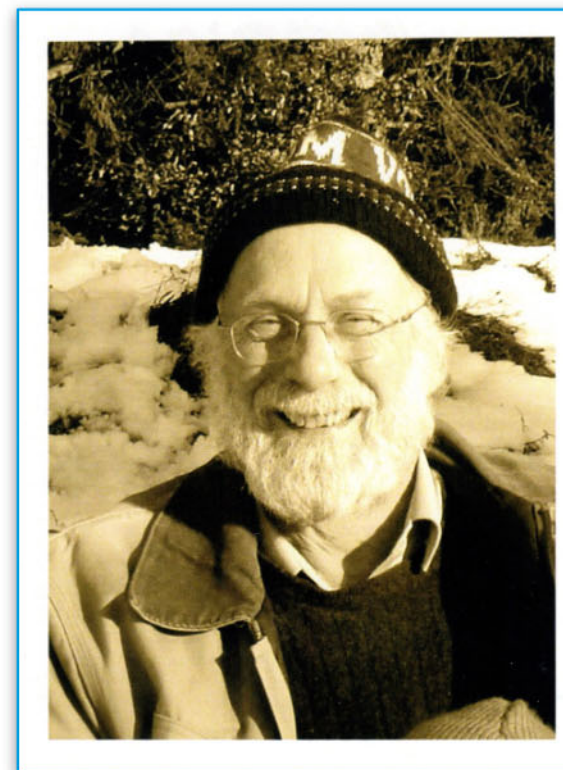


Introduction to SNP

The dream come true

THANK YOU FRANCIS !!!

- Polarising neutron beams with Heusler alloy
- Cryoflipper (no wavelength and stray field dependency)
- **Cryogenic Polarisation Analysis Device** for SNP
(with major contributions by S. Pujol and P. J. Brown)
- ^3He spin filter cell and pumping station



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THE EUROPEAN NEUTRON SOURCE


NEUTRONS
FOR SOCIETY



Introduction to SNP

Milestones (by Francis Tasset)

- 1963 Blume and Maleev independently publish the equations
- 1969 Riste, Moon and Koehler analyse polarisation
- 1973 Alperin demonstrates that 3D polarisation analysis is possible
- 1976 FT's sabbatical at ORNL working on superconducting Nb
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- 1998 SNP is finally accepted and adopted by the experts
- 2002 Cryopad III (first project of ILL Millennium Programme)

$$\begin{aligned} \mathbf{P}'I &= \mathbf{P}(N^2 - \mathbf{M}_\perp \mathbf{M}_\perp^*) \\ &+ 2\Re[\mathbf{M}_\perp(\mathbf{P} \cdot \mathbf{M}_\perp^*)] \\ &+ 2\Re(N^* \mathbf{M}_\perp) \\ &+ \mathbf{P} \times 2\Im(N^* \mathbf{M}_\perp) \\ &- \Im(\mathbf{M}_\perp \times \mathbf{M}_\perp^*) \end{aligned}$$

$$\begin{aligned} I &= N^2 + \mathbf{M}_\perp \mathbf{M}_\perp^* \\ &+ 2\Re(\mathbf{P} \cdot N^* \mathbf{M}_\perp) \\ &+ \mathbf{P} \cdot \Im(\mathbf{M}_\perp \times \mathbf{M}_\perp^*) \end{aligned}$$

Blume (1963) Phys. Rev. **130** 1670

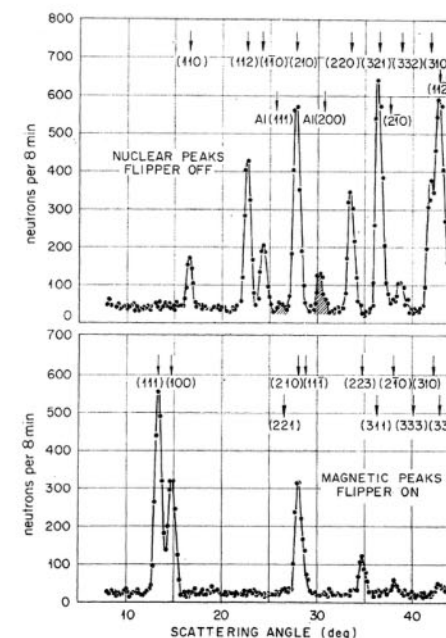
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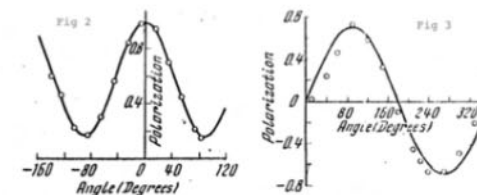
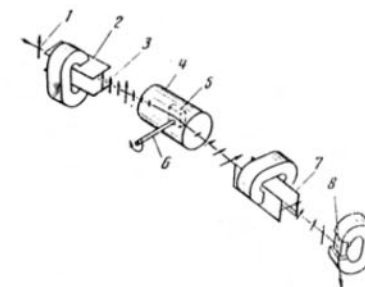
Moon et al. (1969) Phys. Rev. **181** 920



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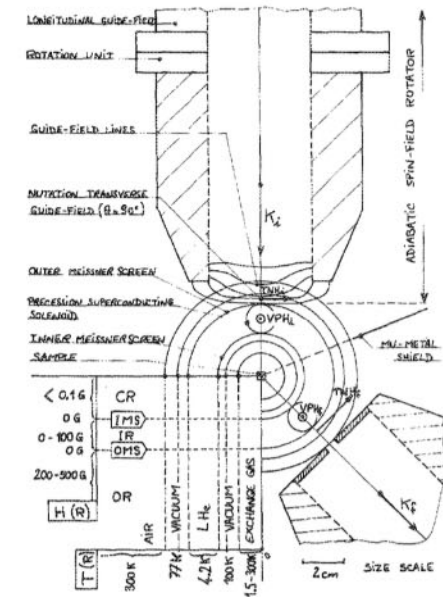
Alperin (1973) International Conference on Magnetism, Moscow, ed., Proc. ICM-73



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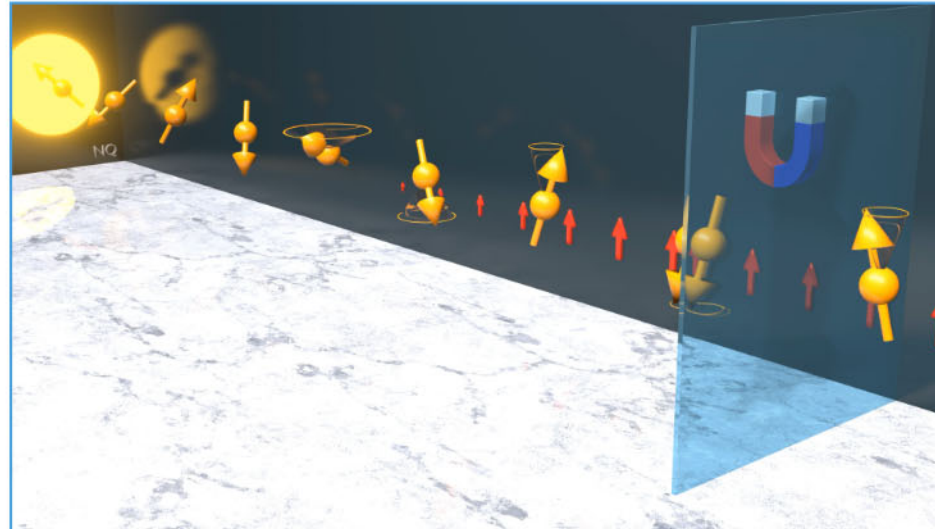


F.Tasset et al. (1988) J. Appl. Phys. **63** 3606



The Cryopad

Manipulation of polarised neutron beams



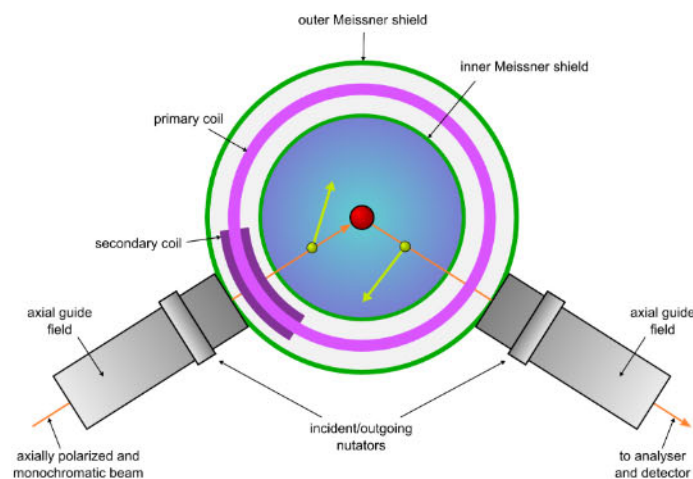
Francis Tasset's polarisation lab



The Cryopad

Manipulation of polarised neutron beams

top view



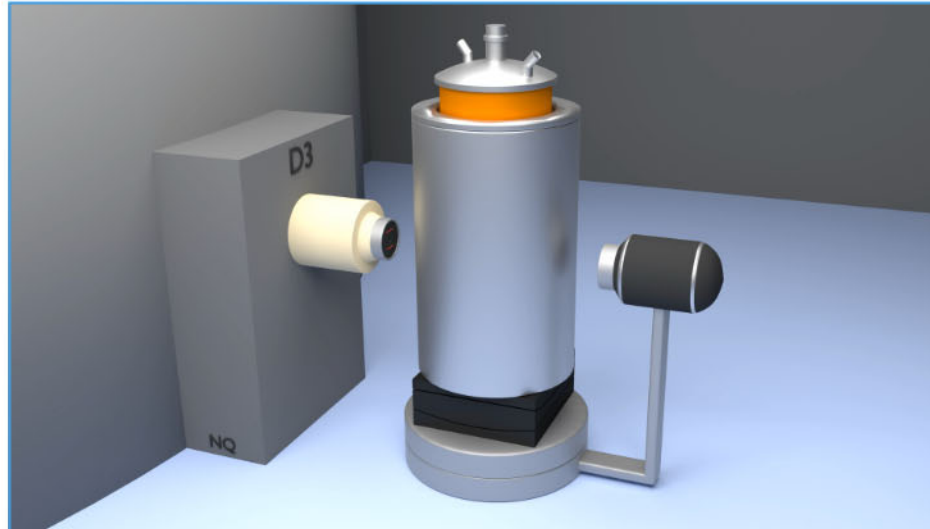
mounted on D3





The Cryopad

Making magic happen



Local coordination system:

$$\mathbf{x} \parallel \mathbf{Q}$$

\mathbf{z} vertical

\mathbf{y} completes right-handed set

$$\mathbf{M}_{\perp,x} = 0$$

$$\mathbf{P}_i \parallel x \quad \mathbf{P}_f \parallel y$$

$$P_{f,i} = P_{yx} = \frac{n^+ - n^-}{n^+ + n^-}$$



The Cryopad

The polarisation matrix

Final neutron polarisation is related to initial one by

$$\mathbf{P}_f = \mathcal{P} \mathbf{P}_i + \mathbf{P}'$$

rotational part
created/annihilated polarisation

$$\mathbf{P}_{f,i} = \begin{pmatrix} p_{fx} \frac{p_{ix}(N^2 - M_{\perp}^2) - J_{yz}}{I_x} & p_{fx} \frac{-p_{iy}J_{nz} - J_{yz}}{I_y} & p_{fx} \frac{p_{iz}J_{ny} - J_{yz}}{I_z} \\ p_{fy} \frac{p_{ix}J_{nz} + R_{ny}}{I_x} & p_{fy} \frac{p_{iy}(N^2 + M_{\perp y}^2 - M_{\perp z}^2) + R_{ny}}{I_y} & p_{fy} \frac{p_{iz}R_{yz} + R_{ny}}{I_z} \\ p_{fz} \frac{-p_{ix}J_{ny} + R_{nz}}{I_x} & p_{fz} \frac{p_{iy}R_{yz} + R_{nz}}{I_y} & p_{fz} \frac{p_{iz}(N^2 - M_{\perp y}^2 + M_{\perp z}^2) + R_{nz}}{I_z} \end{pmatrix}$$

Easily distinguish **nuclear**, **magnetic**, **chiral** and **interference** terms.

$$\begin{aligned} I_x &= N^2 + M_{\perp}^2 + p_{ix}J_{yz} & R_{yz} &= 2\Re(M_{\perp y}M_{\perp z}^*) & R_{nz} &= 2\Re(NM_{\perp z}^*) \\ I_y &= N^2 + M_{\perp}^2 + p_{iy}R_{ny} & J_{yz} &= 2\Im(M_{\perp y}M_{\perp z}^*) & J_{ny} &= 2\Im(NM_{\perp y}^*) \\ I_z &= N^2 + M_{\perp}^2 + p_{iz}R_{nz} \end{aligned}$$



The Cryopad

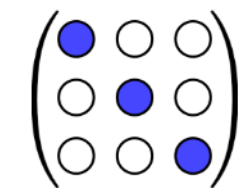
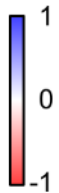
The polarisation matrix

Final neutron polarisation is related to initial one by

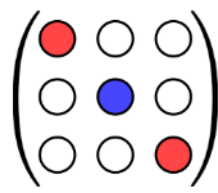
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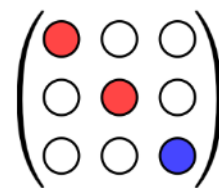
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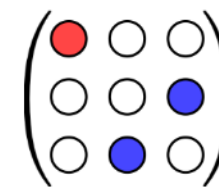
nuclear



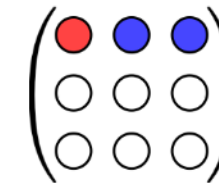
purely magnetic
 $M_{\perp z} = 0$



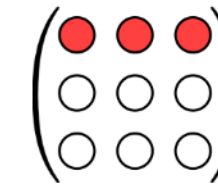
purely magnetic
 $M_{\perp y} = 0$



purely magnetic
 $M_{\perp y} = M_{\perp z}$



left-handed helix
 $iM_{\perp y} = M_{\perp z}$



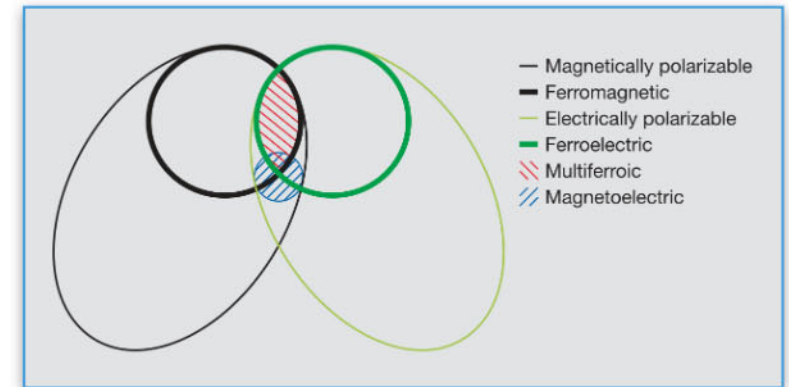
right-handed helix
 $iM_{\perp y} = -M_{\perp z}$



Examples

Magnetoelectric and multiferroic materials

- Magnetoelectric effect predicted by P. Curie in 1894
- Crosstalk between magnetisation and electric polarisation
- Multiferroics possess at least two (anti)ferroic orders
- Used as magnetic field sensors and transducers
- Immense potential for non-volatile data storage applications (faster and consumes less energy)
- spin-transfer torque: ~ 10 fJ
- capacitive ME device: ~ 1 aJ (10^4 less)



Eerenstein et al. (2006) Nature **442** 759

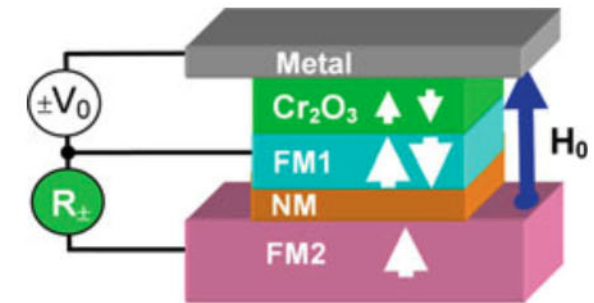


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Magnetic Random Access Memory



Kleemann and Binek (2013)
Multiferroic and magnetoelectric materials

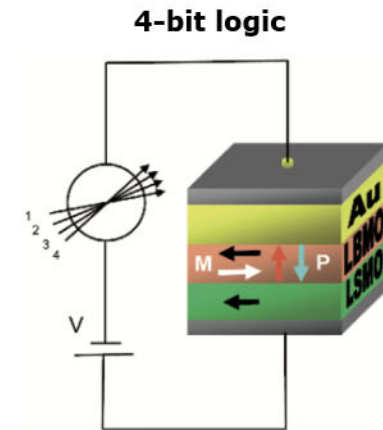


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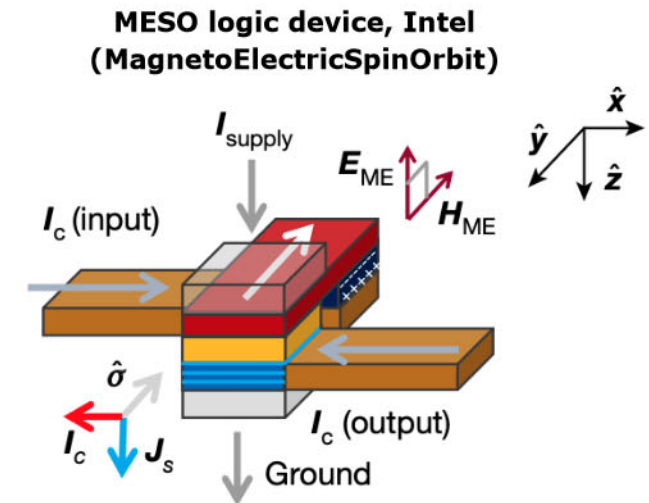
Kleemann and Borisov (2013)
Multiferroic and magnetoelectric materials for spintronics



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Manipatruni et al. (2019) Nature **565** 35



Examples

Cr₂O₃ - the prototype magnetoelectric

- Predicted by Dzyaloshinskii in 1960
- First observation of linear magnetoelectric effect
- $P_i = \alpha_{ij}H_j$ or $M_i = \alpha_{ji}E_j$
- Corundum structure, space group $R\bar{3}c$, $T_N = 307$ K
- Magnetic symmetry ($\bar{3}'m'$) breaks spatial and time-inversion

Astrov (1960) Sov. Phys. JETP **11** 708

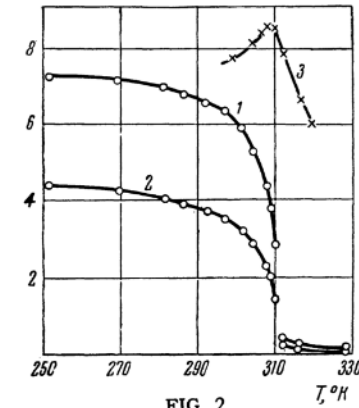
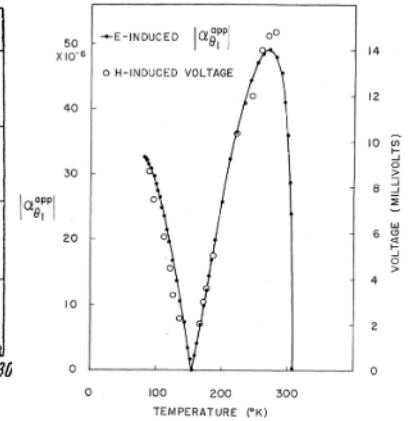


FIG. 2



Rado and Folen (1961) Phys. Rev. Lett **7** 310

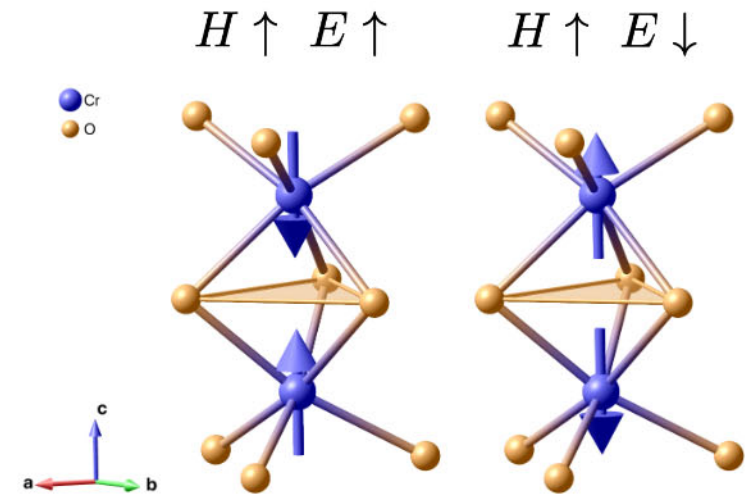


Examples

Cr_2O_3 - sign of magnetoelectric coefficient

- Problem: presence of 180° magnetic domains
- Annealing under field achieves domain selection
- Domains indistinguishable with unpolarised neutrons
- SNP can tell the difference!
- b vertical, $(-1\ 0\ 2)$: $P_i \parallel z$ rotates towards $\pm x$
- information about sign of magnetoelectric coefficient

Brown, Forsyth and Tasset (1998) J. Phys.: Condens. Matter. **10** 663





Examples

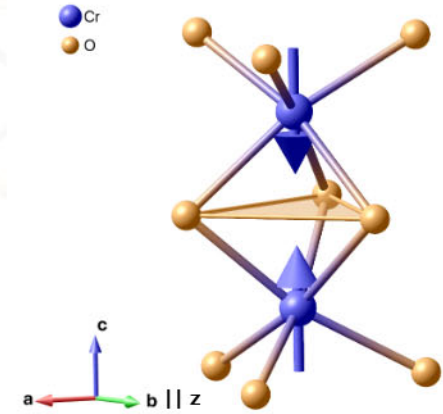
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- Domains indistinguishable with unpolarised neutrons
- SNP can be different for different domains
- $b \parallel z$, moments along $c \rightarrow$ no z component
- b vertical
- $(-1 \ 0 \ 2)$ at room temperature: $N \approx M_{\perp y}$
- $\Im(N) = 0 \quad \Re(M_{\perp}) = 0 \rightarrow$ No real interference terms

Brown, Forsyth and Tassel (1998) J. Phys.: Condens. Matter, 10 663

$$R_{ny} = 2\Re(NM_{\perp y}^*) \quad J_{ny} = 2\Im(NM_{\perp y}^*)$$



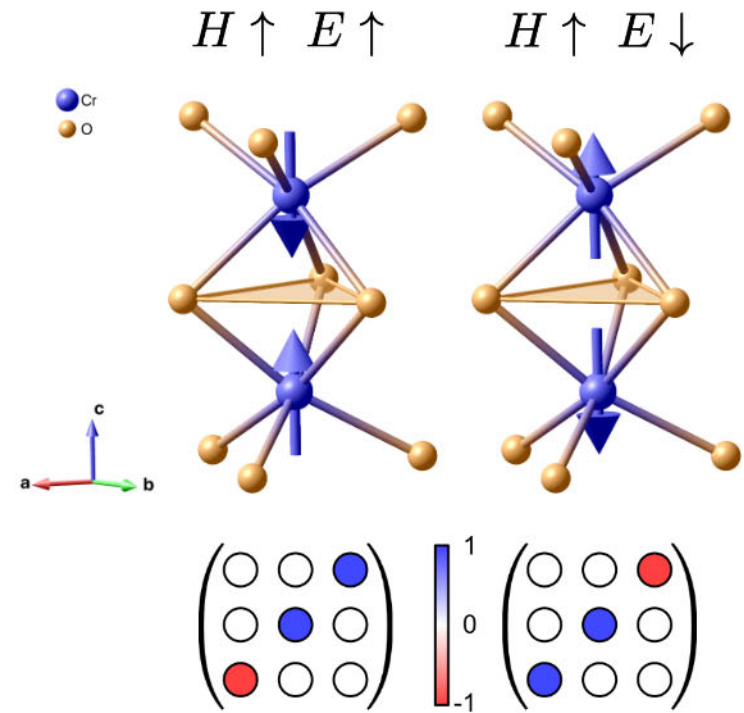


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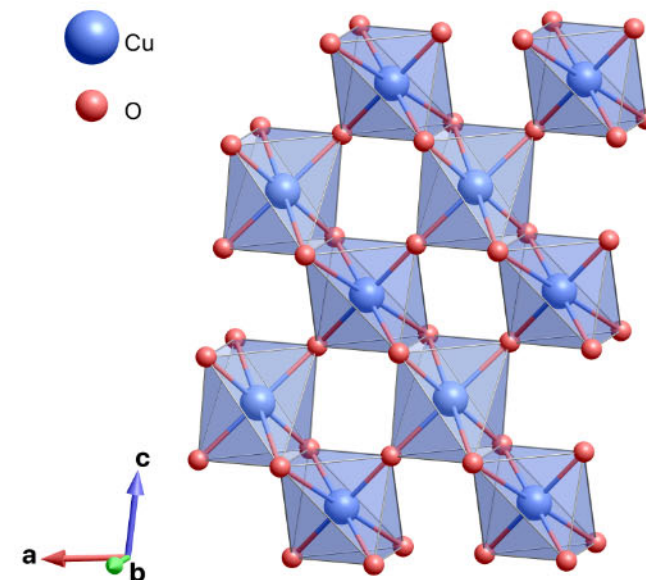




Examples

CuO - a high-T multiferroic

- so far the only known binary multiferroic compound
- monoclinic space group $C2/c$
- spiral magnetic ordering induces electric polarisation at 230 K
- DFT + MC: applying pressure drives multiferroic state towards RT
- recently confirmed by neutron diffraction under pressure

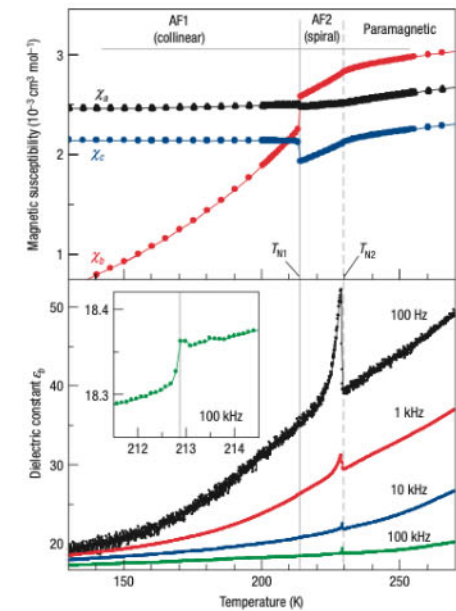




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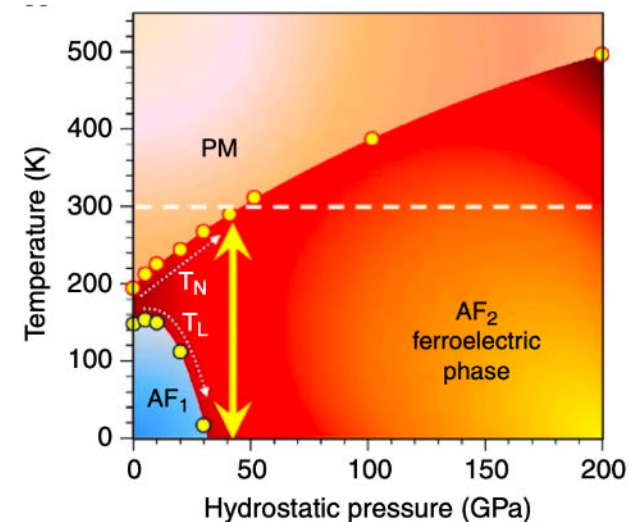
Kimura et al. (2008) *Nat. Mater.* **7** 291



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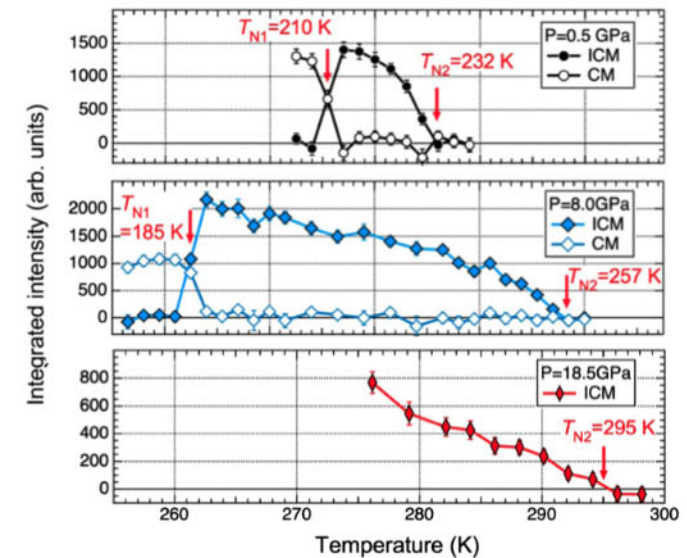
Rocquefelte et al. (2013) *Nat. Mater.* **4** 2511



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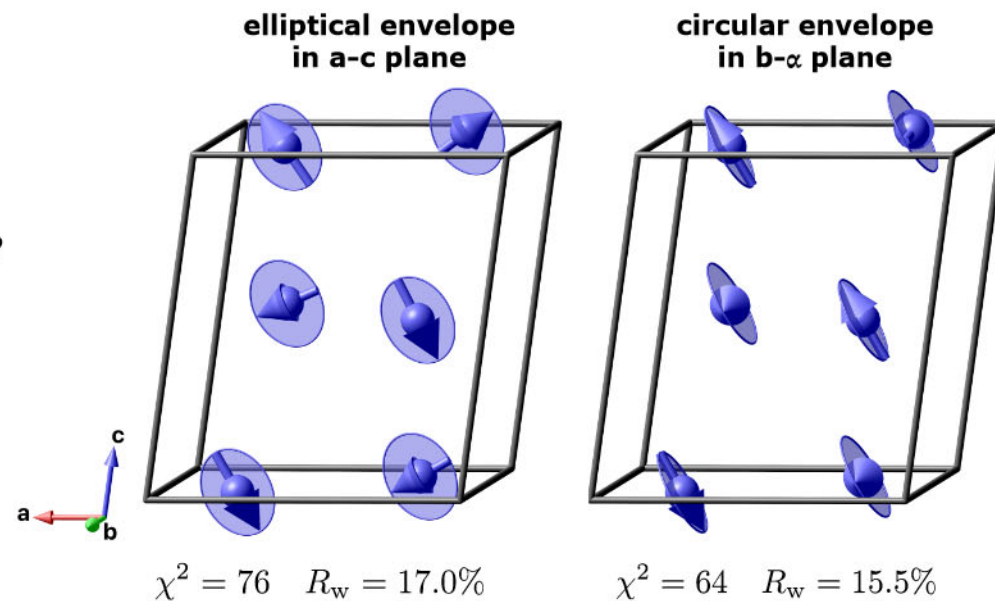
Terada et al. (2022) *Phys. Rev. Lett.* **129** 217601



Examples

CuO - a high-T multiferroic

- CuO is a building block of high-T_c cuprates
- magnetism in CuO → hole pairing mechanism?
- unpolarised neutrons: model ambiguity



Forsyth, Brown et al. (1988) J. Phys. C: Solid State Phys **21** 2917

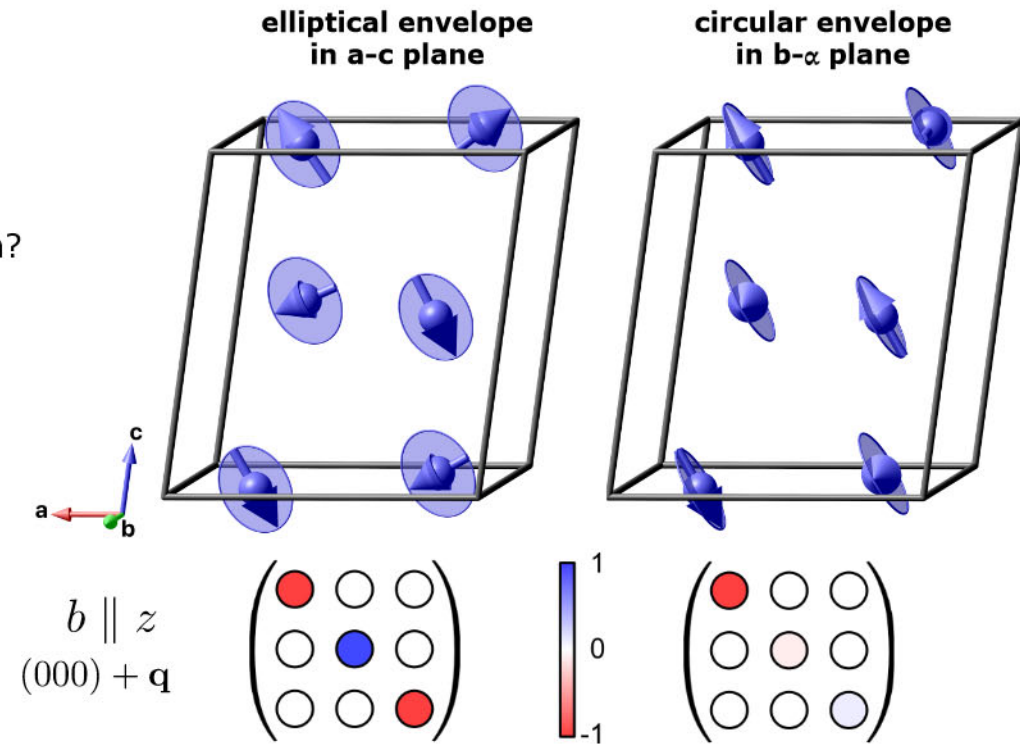
Brown, ..., Forsyth, ... and Tasset (1991) J. Phys.: Condens. Matter. **3** 4281



Examples

CuO - a high-T multiferroic

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- unpolarised neutrons: model ambiguity
- correct model determined by SNP
- superior to unpolarised neutrons especially for weak reflections and non-collinear structures

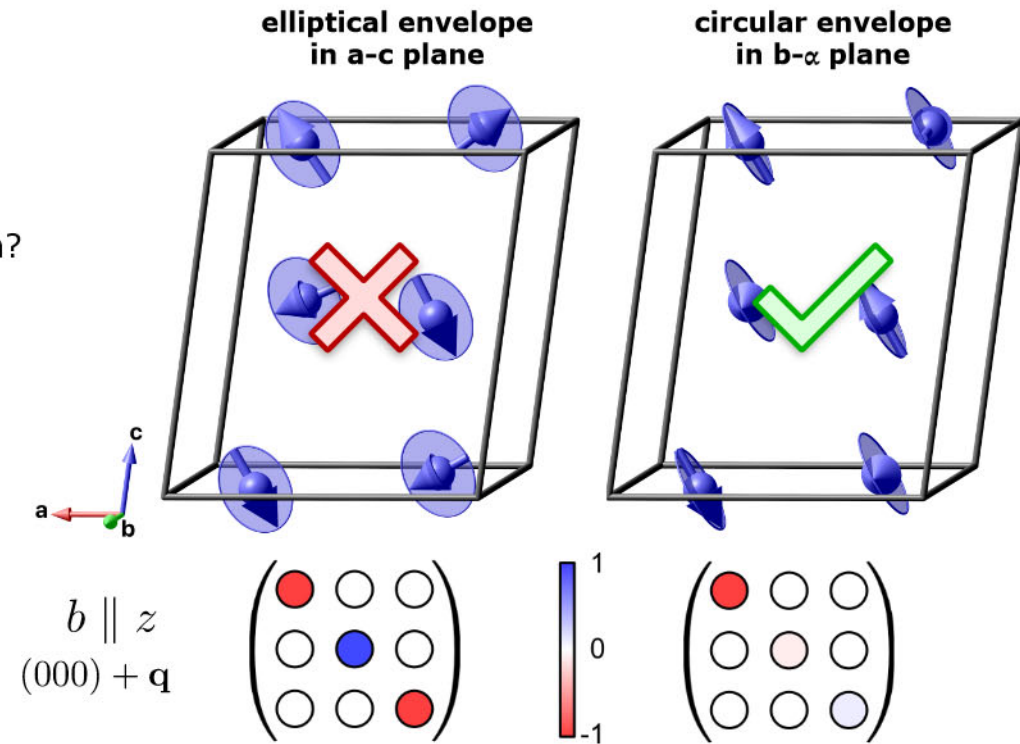




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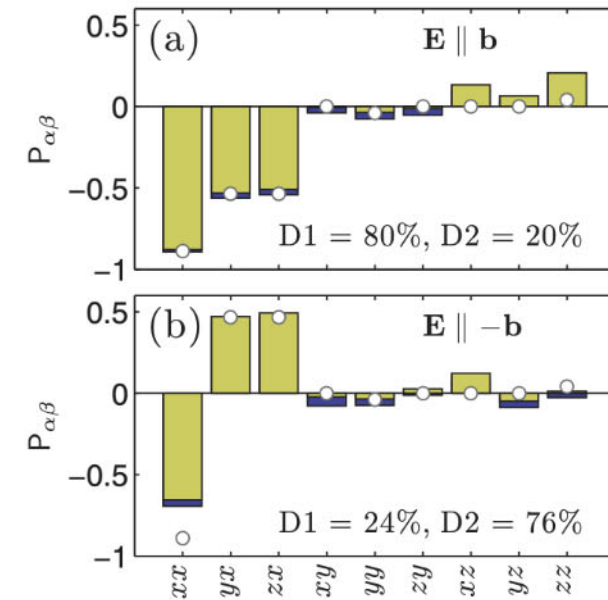




Examples

CuO - switching behaviour

- strong coupling between magnetic order and electric polarisation
- electric field control of chiral domains
- magnetic chirality is inaccessible with unpolarised neutrons
- SNP can tell the difference



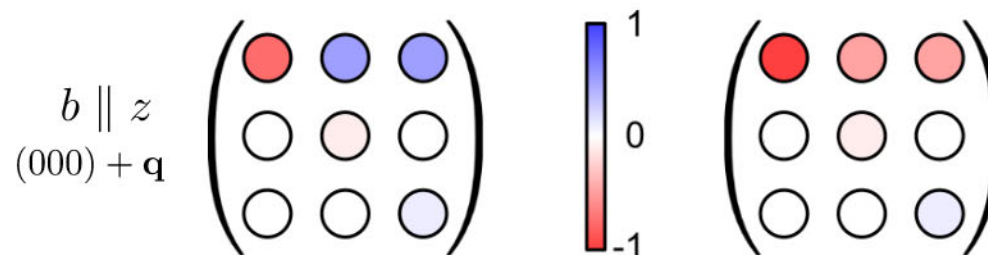
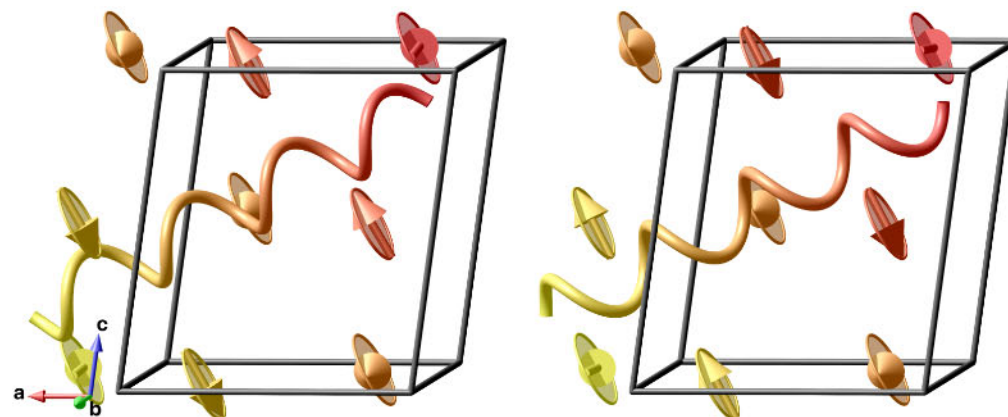
Babkevich et al. (2012) Phys. Rev.B **85** 134428



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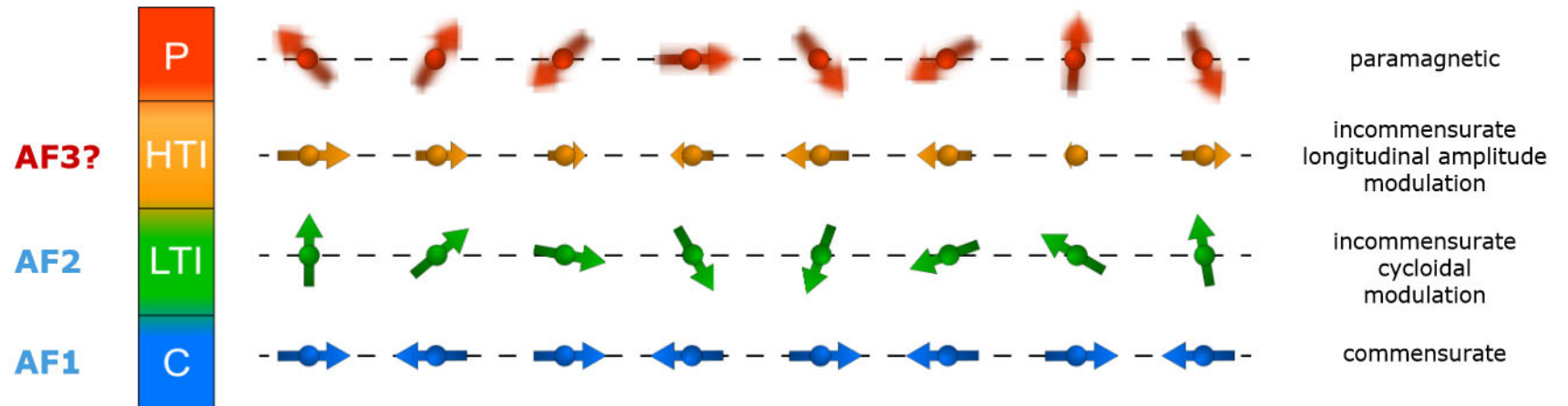




Examples

CuO - search for the AF3 phase

usual sequence of magnetic phases in multiferroics (e.g. TbMnO_3 , $\text{Ni}_3\text{V}_2\text{O}_8$):

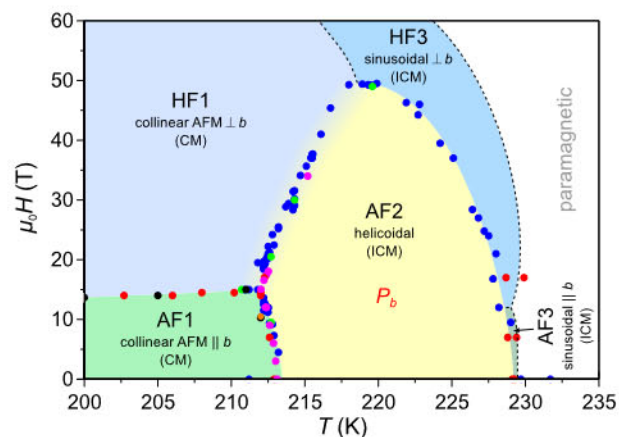




Examples

CuO - search for the AF3 phase

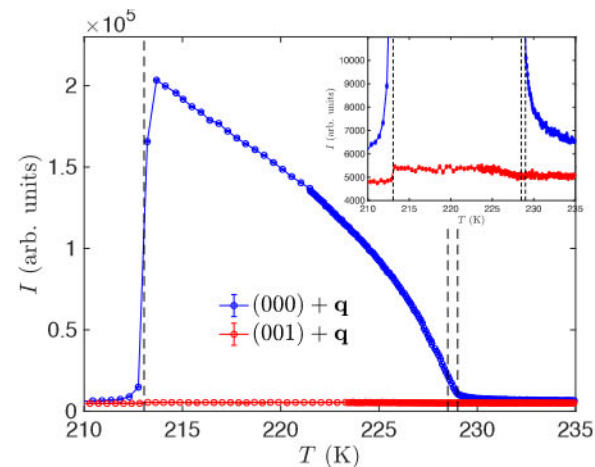
macroscopic methods



0.5 K stability range!

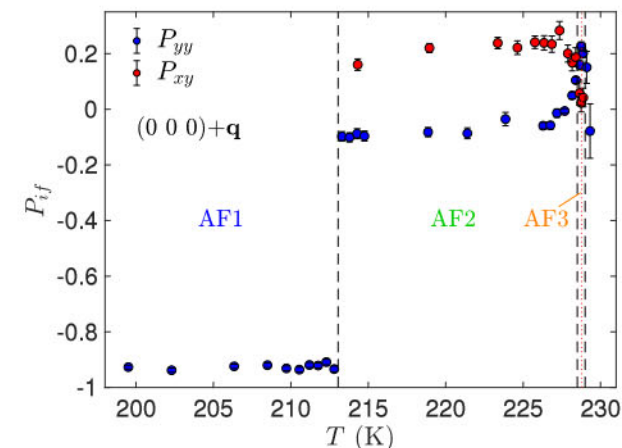
Wang, Qureshi et al. (2016) *Nat. Commun.* **7** 10295

unpolarised neutrons



No anomalies in I(T)

SNP on P_{yy} and P_{xy} elements



Microscopic proof of AF3!

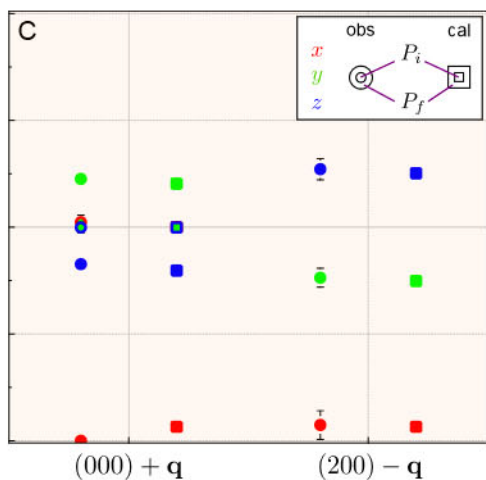
Qureshi et al. (2020) *Sci. Adv.* **6** eaay7661



Examples

CuO - search for the AF3 phase

observed vs. calculated

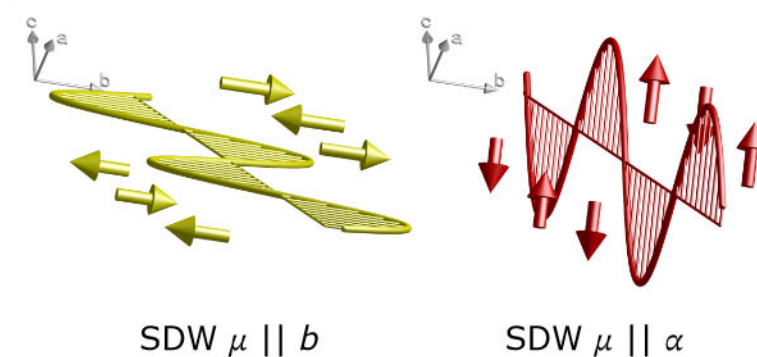


analysed with



Qureshi (2019) J. Appl. Cryst. **52** 175

phase coexistence



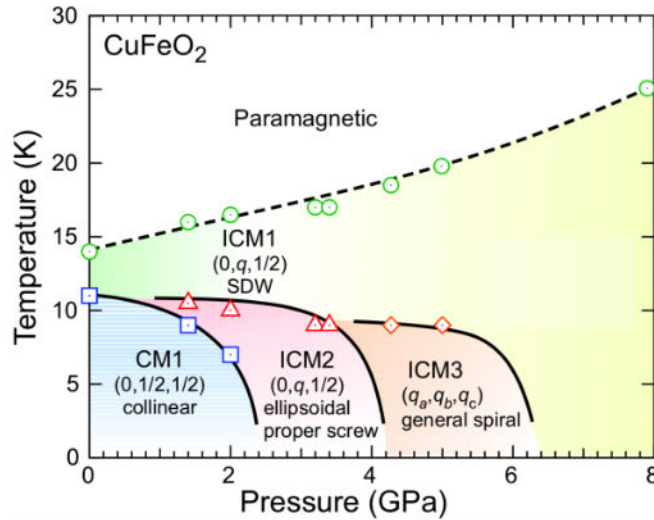


Examples

CuFeO₂ - pressure-induced multiferroicity

Complex p-T phase diagram

including collinear spin-density waves and non-collinear spiral structures

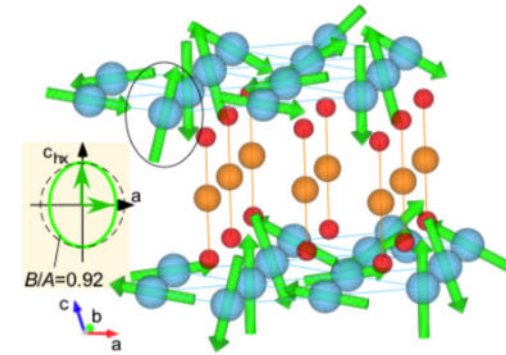


Lack of information

Limited number of observations (integrated intensities) did not allow to reveal:

Terada et al. (2014) Phys. Rev. B **83** 220403(R)

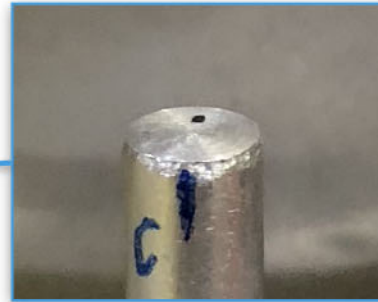
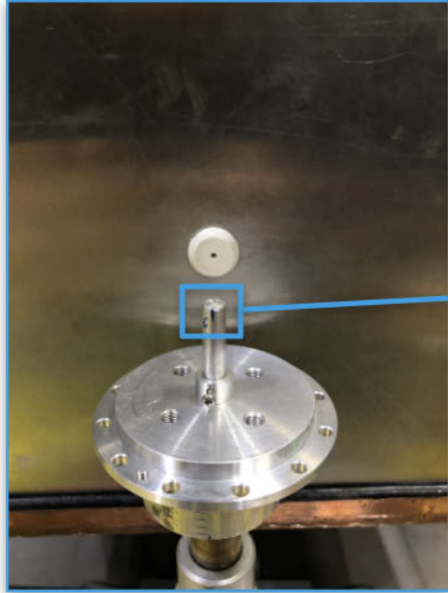
- stability region of ICM1
- ellipsoidicity ratio of ICM2
- magnetic structure of ICM3





Examples

CuFeO₂ - pressure-induced multiferroicity



0.5 x 0.5 x 0.2 mm³



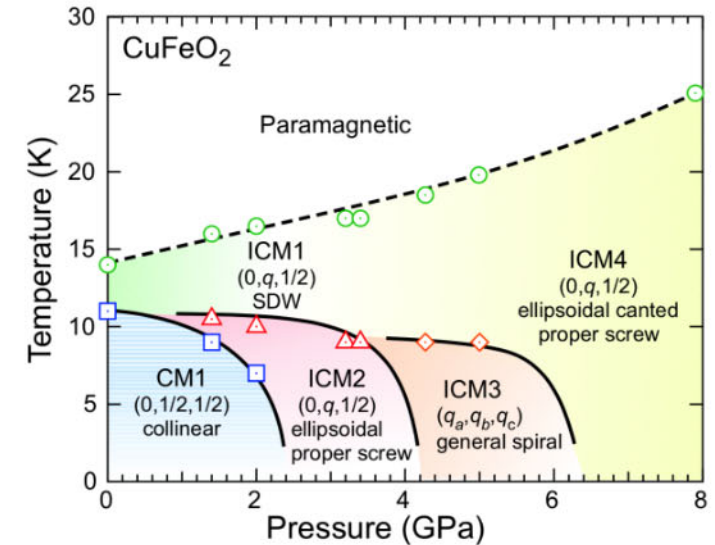
non-magnetic pressure cell (up to 10 GPa)



Examples

CuFeO₂ - pressure-induced multiferroicity

- SNP under pressure: It's possible!!!
- Determined 5 complex magnetic structures at 3 different pressures in 7 days with a 0.05 mm³ sample
- SNP is an excellent approach to study novel pressure-induced phenomena associated with complex magnetic order



Terada, Qureshi et al. (2018) Nat. Commun. **9** 4368



Conclusions

SNP and the future

- SNP is a powerful technique without alternatives for certain scientific challenges
- Magnetic domain sensitivity makes it the ideal probe for switching behaviour of smart materials
- Weak signals (small moments, small samples, thin films) are no problem as long as different from background
- Importance for future spintronic candidates, but also emerging field of twistrionics (2D materials, few-layer limit)



Conclusions

SNP and the future

- SNP is a powerful technique without alternatives for certain scientific challenges
- Magnetic domain sensitivity makes it the ideal probe for switching behaviour of smart materials
- Weak signals (small moments, small samples, thin films) are no problem as long as different from background
- Importance for future spintronic candidates, but also emerging field of twistronics (2D materials, few-layer limit)
- D3 at 50%: overload of 5.9 (!) in College 5B in current proposal round
- We need D3 at 100% in order not to miss out on excellent science!



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